

# Photons at rest around boson stars

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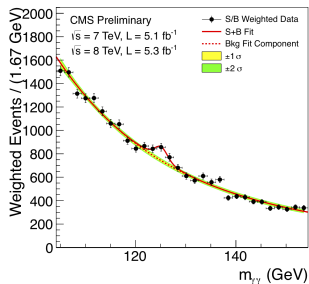
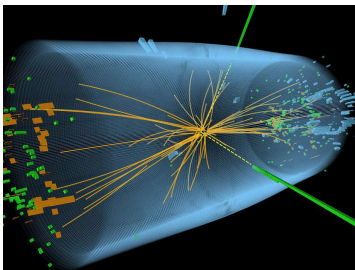
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# What is a boson star ?

- Localized configuration of a self-gravitating complex scalar field.
- Introduced in the 1960s by Bonazzola, Pacini, Kaup and Ruffini.
- Corresponds to spin-0 particle  $\implies$  *boson*.
- At least one scalar field in nature : *Higgs boson*.



# Boson star model

- Scalar field has a  $U(1)$  symmetry:

$$\Phi \longrightarrow \Phi \exp(i\alpha).$$

- The Lagrangian of the field is

$$\mathcal{L}_M = -\frac{1}{2} [g^{\mu\nu} \nabla_\mu \bar{\Phi} \nabla_\nu \Phi + V(|\Phi|^2)].$$

$V$  is a potential (for a free field  $V = m^2/\hbar^2 |\Phi|^2$ )

- The Lagrangian of gravity is

$$\mathcal{L}_g = \frac{1}{16\pi} R$$

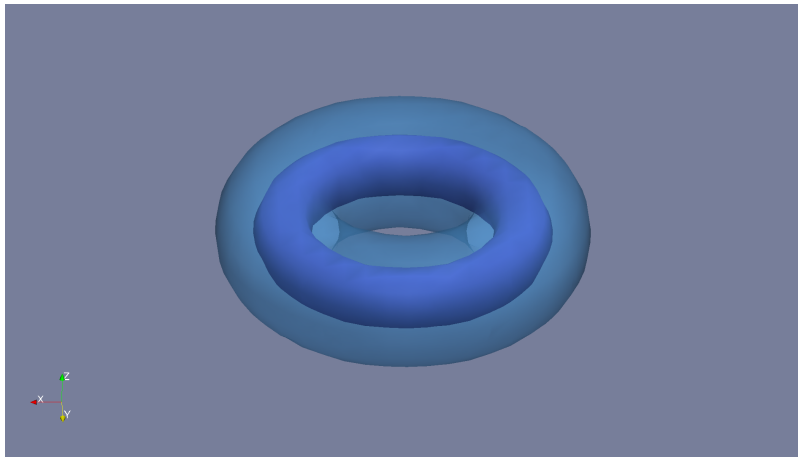
**The variation of the equation gives the Einstein-Klein-Gordon system.**

## Ansatz for the field

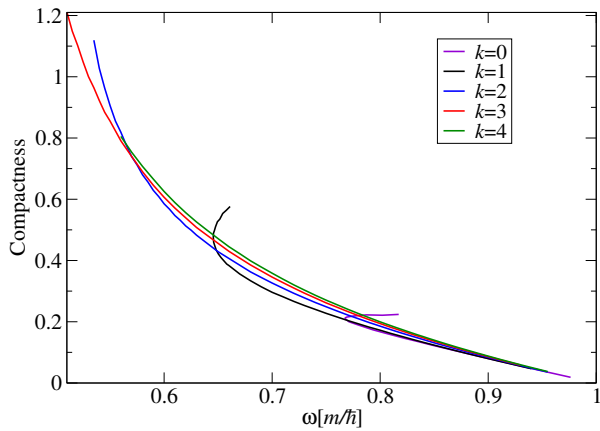
$$\Phi = \phi \exp [i (\omega t - k\varphi)]$$

- $U(1) \implies$  the action does not depend on  $(t, \varphi)$ .
- $\phi$  and the metric fields depend only on  $(r, \theta)$ .
- $k$  and  $\omega$  appear as parameters in the equations.
- $k$  is an integer and  $\omega$  a real number.
- $k = 0$  corresponds to spherically symmetric configurations.

# Field configuration ; $k = 2$



# Compactness



# Effective potential method

- Metric in quasi-isotropic coordinates :

$$g_{\mu\nu} dx^\mu dx^\nu = -N^2 dt^2 + A^2 (dr^2 + r^2 d\theta^2) + B^2 r^2 \sin^2 \theta (d\varphi + \beta^\varphi dt)^2$$

- Light rings : closed circular orbits of photons.

$$U^\alpha = (U^t, U^r, 0, U^\varphi)$$

- Two conserved quantities  $U_\alpha (\partial_t)^\alpha = -E$  and  $U_\alpha (\partial_\varphi)^\alpha = L$ .
- Null geodesics  $U_\alpha U^\alpha = 0$  leads to  $(U^r)^2 + V_{\text{eff}}(r) = 0$

$$V_{\text{eff}} = \frac{1}{A^2} \left[ -\frac{(\beta^\varphi L + E)^2}{N^2} + \frac{L^2}{B^2 r^2} \right]$$

# Circular orbits

- Circular orbits require  $V_{\text{eff}} = 0$  and  $\partial_r V_{\text{eff}} = 0$
- First condition implies

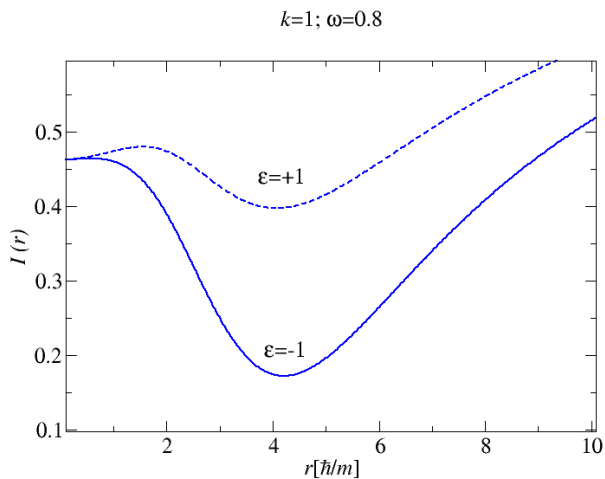
$$\frac{E}{L} = -\beta^\varphi + \epsilon \frac{N}{Br} \quad ; \quad \epsilon = \pm 1$$

- Contrary to the massive case only the ratio  $E/L$  is constrained.
- Second condition reduces to finding the zeros of

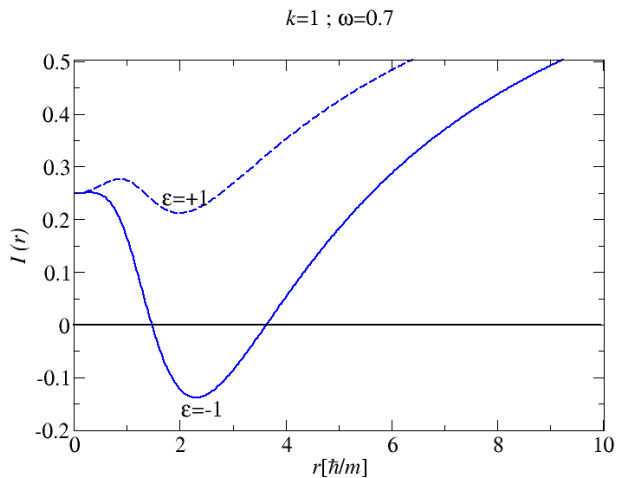
$$I(r) = \left( \epsilon \frac{\partial_r \beta^\varphi}{NB} \right) r^2 + \left( \frac{\partial_r B}{B^3} - \frac{\partial_r N}{NB^2} \right) r + \frac{1}{B^2}$$



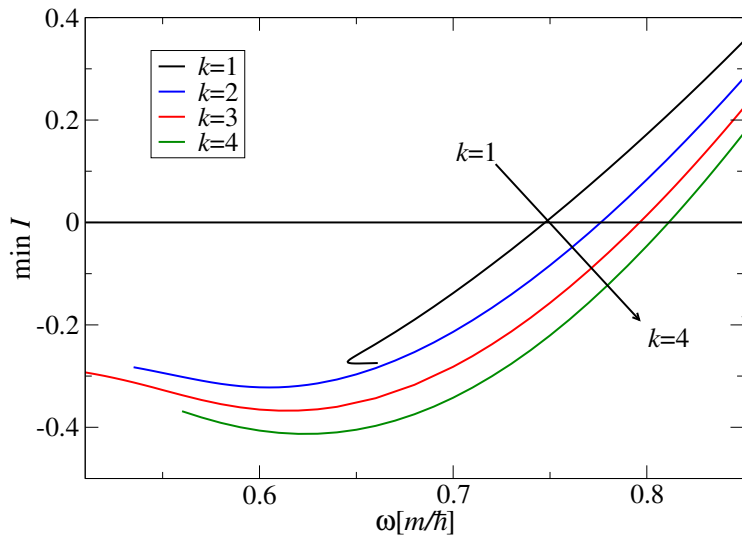
# $I(r)$ without light rings



# $I(r)$ with light rings



# Existence of light rings



# Properties

The orbital frequency is given by  $\frac{d\varphi}{dt} = -\frac{N}{Br} - \beta\varphi$ .

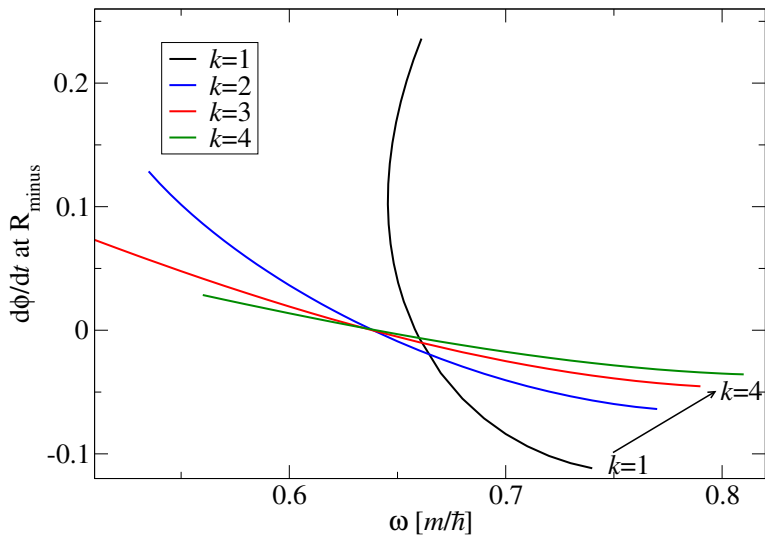
## Boson star case

- Two light rights for relativistic enough configurations.
- The outer one is unstable and retrograde.
- The outer one is stable (and changes type).

## Black hole case

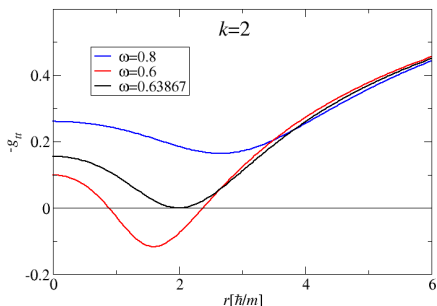
- Two light rings for Kerr black holes.
- One prograde and one retrograde.
- Both unstable.

# Orbital frequency of the inner light ring



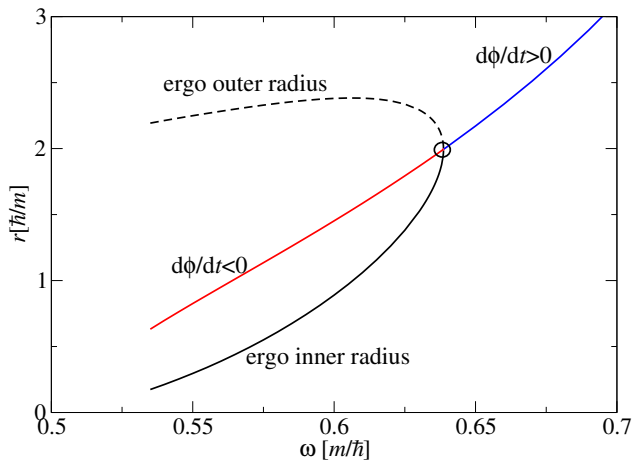
# Light points

- On each sequence : a photon **at rest** :  $U^\mu = (1, 0, 0, 0)$ .
- Null condition :  $U_\mu U^\mu = 0 \implies g_{tt} = 0$ . The light point lies exactly on the boundary of an ergoregion.
- Additional condition : geodesic equation  $\implies \partial_r g_{tt} = 0$

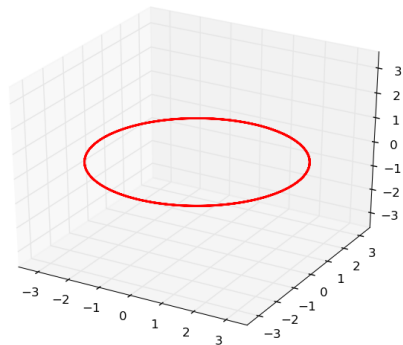


# Light points and ergoregions

The light points lie where an ergoregion just starts to develop



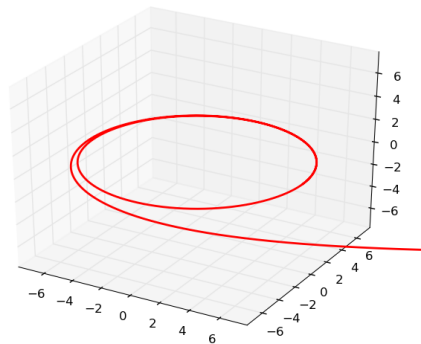
# Inner light ring



Direct integration with Gyoto ;  $k = 2$  and  $\omega = 0.7$ .

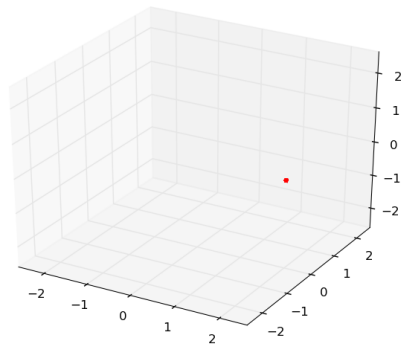


# Outer light ring



Direct integration with Gyoto ;  $k = 2$  and  $\omega = 0.7$ .

# Light point



Direct integration with Gyoto ;  $k = 2$  and  $\omega = 0.6387$ .

# Last words

- Boson stars are indeed very compact.
- Existence of two light rings (one stable).
- On each sequence existence of a single light point, where the photon is at rest.
- Not present in the BH case. Schwarzschild admit an unstable photon at rest trajectory, exactly on the horizon.
- Boson stars are good testbeds for strong relativistic effects.